

Efficient Optimal Control for Robust Quantum Gates in Liouville Space

Michael Goerz

Universität Kassel, Germany

April 24, 2014

QUAINT Workshop 2014
Sandbjerg Estate, Aarhus University, Denmark

The Problem of Robustness

Problem

Problem of real world quantum engineering: Robustness

- Robustness with respect to **decoherence**
⇒ Liouville space
- Robustness with respect to **experimental fluctuations and uncertainties**



The Problem of Robustness

Problem

Problem of real world quantum engineering: Robustness

- Robustness with respect to **decoherence**
⇒ Liouville space
- Robustness with respect to **experimental fluctuations and uncertainties**



Analytical solutions? Probably not, in most cases!

The Problem of Robustness

Problem

Problem of real world quantum engineering: Robustness

- Robustness with respect to **decoherence**
⇒ Liouville space
- Robustness with respect to **experimental fluctuations and uncertainties**

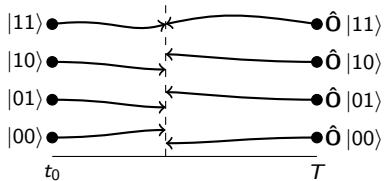
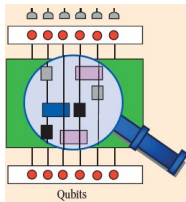


Analytical solutions? Probably not, in most cases!

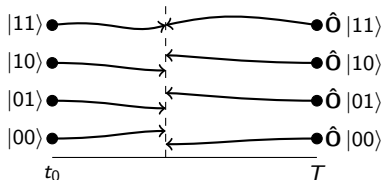
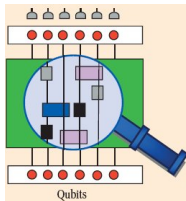
⇒ Optimal Control

- find controls at the quantum speed limit (avoid decoherence)
- ask for robustness *explicitly* in the optimization
 - Optimize in Liouville space... efficiently?
 - Optimize over ensembles of Hamiltonians

Gate Optimization for Open Quantum Systems



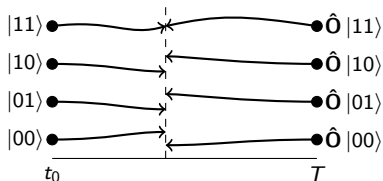
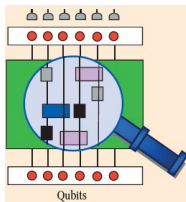
Gate Optimization for Open Quantum Systems



Lift $J_T = 1 - \frac{1}{d} \sum_{i=1}^d \Re e \langle \Psi_i | \hat{O}^\dagger \hat{U}(T, 0, \epsilon) | \Psi_i \rangle$ to Liouville space.

Kallush & Kosloff, PRA 73, 032324 (2006),
 Schulte-Herbrüggen et al., JPB 44, 154013 (2011),
 Ohtsuki, NJP 12, 045002 (2010), ...

Gate Optimization for Open Quantum Systems

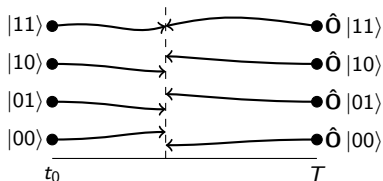
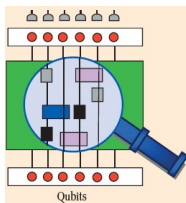


Lift $J_T = 1 - \frac{1}{d} \sum_{i=1}^d \Re \langle \Psi_i | \hat{O}^\dagger \hat{U}(T, 0, \epsilon) | \Psi_i \rangle$ to Liouville space.

Kallush & Kosloff, PRA 73, 032324 (2006),
 Schulte-Herbrüggen et al., JPB 44, 154013 (2011),
 Ohtsuki, NJP 12, 045002 (2010), ...

$$\Rightarrow J_t = 1 - \frac{1}{d^2} \sum_{j=1}^{d^2} \Re \left[\text{tr} \left[\hat{O} \hat{\rho}_j(0) \hat{O}^\dagger \hat{\rho}_j(T) \right] \right]$$

Gate Optimization for Open Quantum Systems



Lift $J_T = 1 - \frac{1}{d} \sum_{i=1}^d \Re \langle \psi_i | \hat{\mathbf{O}}^\dagger \hat{\mathbf{U}}(T, 0, \epsilon) | \psi_i \rangle$ to Liouville space.

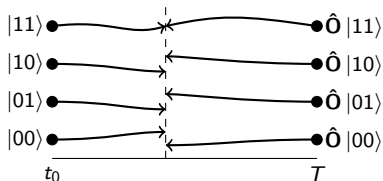
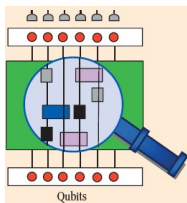
Kallush & Kosloff, PRA 73, 032324 (2006),
Schulte-Herbrüggen et al., JPB 44, 154013 (2011),
Ohtsuki, NJP 12, 045002 (2010), ...

$$\Rightarrow J_t = 1 - \frac{1}{d^2} \sum_{j=1}^{d^2} \Re \left[\text{tr} \left[\hat{\mathbf{O}} \hat{\rho}_j(0) \hat{\mathbf{O}}^\dagger \hat{\rho}_j(T) \right] \right]$$

$$\hat{\rho}_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \quad \hat{\rho}_2 = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \quad \hat{\rho}_3 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \quad \dots$$

d^2 matrices to propagate! (16 for two-qubit gate)

Gate Optimization for Open Quantum Systems



Lift $J_T = 1 - \frac{1}{d} \sum_{i=1}^d \Re \langle \psi_i | \hat{O}^\dagger \hat{U}(T, 0, \epsilon) | \psi_i \rangle$ to Liouville space.

Kallush & Kosloff, PRA 73, 032324 (2006),
Schulte-Herbrüggen et al., JPB 44, 154013 (2011),
Ohtsuki, NJP 12, 045002 (2010), ...

$$\Rightarrow J_t = 1 - \frac{1}{d^2} \sum_{j=1}^{d^2} \Re \left[\text{tr} \left[\hat{O} \hat{\rho}_j(0) \hat{O}^\dagger \hat{\rho}_j(T) \right] \right]$$

Claim

We only need to propagate **three** matrices (independent of d), instead of d^2 .

A Reduced Set of Density Matrices

No need to characterize the full dynamical map! – much less information required to assess how well a desired unitary is implemented

A Reduced Set of Density Matrices

No need to characterize the full dynamical map! – much less information required to assess how well a desired unitary is implemented

① **Do we stay in the logical subspace?**

No need to characterize the full dynamical map! – much less information required to assess how well a desired unitary is implemented

① Do we stay in the logical subspace?

$$\hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

A Reduced Set of Density Matrices

No need to characterize the full dynamical map! – much less information required to assess how well a desired unitary is implemented

- ① Do we stay in the logical subspace?
- ② Are we unitary, and if yes, did we implement the right gate?

$$\hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

No need to characterize the full dynamical map! – much less information required to assess how well a desired unitary is implemented

- ① Do we stay in the logical subspace?
- ② Are we unitary, and if yes, did we implement the right gate?

$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix},$$

$$\hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

A Reduced Set of Density Matrices

No need to characterize the full dynamical map! – much less information required to assess how well a desired unitary is implemented

① Do we stay in the logical subspace?

② Are we unitary, and if yes, did we implement the right gate?

$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix},$$

$$\hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

E.g. $\hat{\mathbf{O}} = \text{diag}(-1, 1, 1, 1)$;

For $\hat{\mathbf{U}} = \text{diag}(e^{i\phi_{00}}, e^{i\phi_{01}}, e^{i\phi_{10}}, e^{i\phi_{11}})$

using just $\hat{\rho}_1$ will not distinguish $\hat{\mathbf{U}}$ from $\hat{\mathbf{O}}$. ($\hat{\mathbf{U}}\hat{\rho}_1\hat{\mathbf{U}}^\dagger = \hat{\mathbf{O}}\hat{\rho}_1\hat{\mathbf{O}}^\dagger = \hat{\rho}_1$)

A Reduced Set of Density Matrices

No need to characterize the full dynamical map! – much less information required to assess how well a desired unitary is implemented

① Do we stay in the logical subspace?

② Are we unitary, and if yes, did we implement the right gate?

$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix}, \hat{\rho}_2 = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}, \hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

E.g. $\hat{\mathbf{O}} = \text{diag}(-1, 1, 1, 1)$;

For $\hat{\mathbf{U}} = \text{diag}(e^{i\phi_{00}}, e^{i\phi_{01}}, e^{i\phi_{10}}, e^{i\phi_{11}})$

using just $\hat{\rho}_1$ will not distinguish $\hat{\mathbf{U}}$ from $\hat{\mathbf{O}}$. ($\hat{\mathbf{U}}\hat{\rho}_1\hat{\mathbf{U}}^\dagger = \hat{\mathbf{O}}\hat{\rho}_1\hat{\mathbf{O}}^\dagger = \hat{\rho}_1$)

Optimization States

$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix}, \quad \hat{\rho}_2 = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}, \quad \hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

populations

phases

subspace

Functional

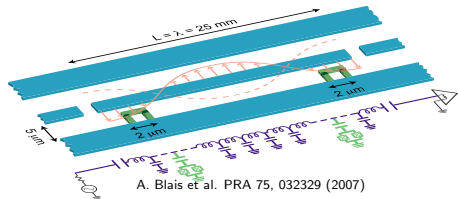
$$J_T = 1 - \sum_{j=1}^3 \frac{w_j}{\text{tr}[\hat{\rho}_j^2(0)]} \Re \left[\text{tr} \left[\hat{\mathbf{O}} \hat{\rho}_j \hat{\mathbf{O}}^\dagger \mathcal{D}[\hat{\rho}_j] \right] \right]$$

- Allow for different weights ($\sum w_j = 1$)
- $J_T = 0$ if and only if $\forall \hat{\rho}_j: \mathcal{D}[\hat{\rho}_j] \equiv \text{target state}$
 \Rightarrow implemented unitary gate $\hat{\mathbf{O}}$.

Example 1

Optimizing for Robustness under Dissipation for a Transmon Gate

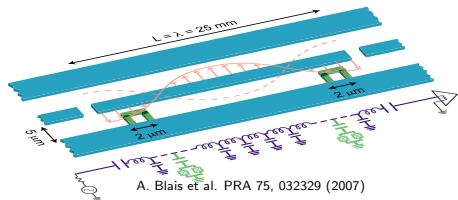
Two Coupled Transmon Qubits



Cavity mediates

- driven excitation of qubit
- interaction between left and right qubit

Two Coupled Transmon Qubits



Cavity mediates

- driven excitation of qubit
- interaction between left and right qubit

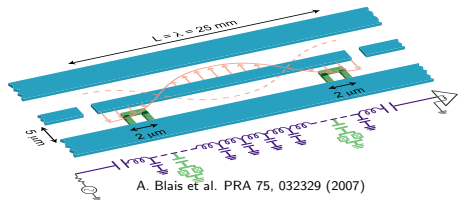
Effective Hamiltonian

$$\hat{H}_{\text{eff}} = \sum_{q=1,2} \sum_{i=0}^{N_q-1} (\omega_i^{(q)} + \chi_i^{(q)}) \hat{n}_i^{(q)} + \sum_{q=1,2} \sum_{i=0}^{N_q-1} g_i^{\text{eff}(q)} \epsilon(t) (\hat{c}_i^{+(q)} + \hat{c}_i^{- (q)})$$

$$+ \sum_{ij} J_{ij}^{\text{eff}} (\hat{c}_i^{- (1)} \hat{c}_j^{+(2)} + \hat{c}_i^{+(1)} \hat{c}_j^{- (2)}).$$

with $\omega_i^{(q)} = i\omega_q - \frac{1}{2}(i^2 - i)\alpha_q$, $\hat{n}_i^{(q)} = |i\rangle\langle i|_q$, $\hat{c}_i^{+(q)} = |i\rangle\langle i-1|_q$

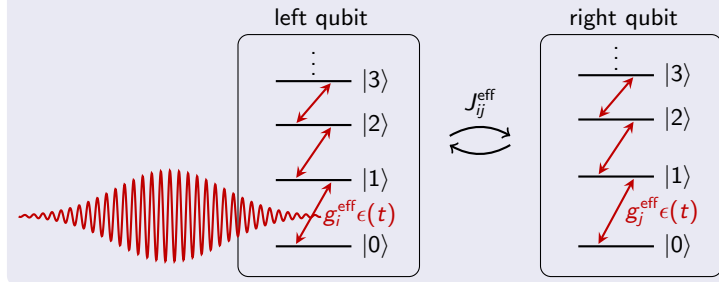
Two Coupled Transmon Qubits



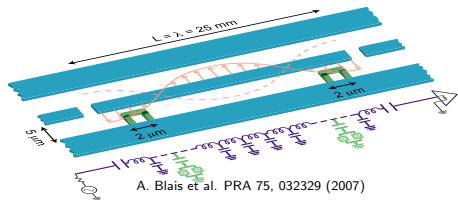
Cavity mediates

- driven excitation of qubit
- interaction between left and right qubit

Effective Hamiltonian



Two Coupled Transmon Qubits



Cavity mediates

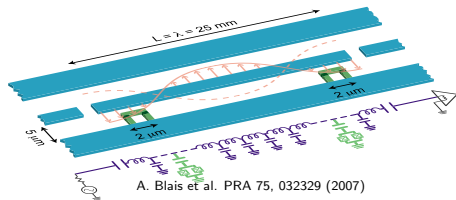
- driven excitation of qubit
- interaction between left and right qubit

Optimization Target

Many gates possible, e.g. \sqrt{i} SWAP:

$$\hat{O} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 \\ 0 & \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

Two Coupled Transmon Qubits



Cavity mediates

- driven excitation of qubit
- interaction between left and right qubit

Master Equation

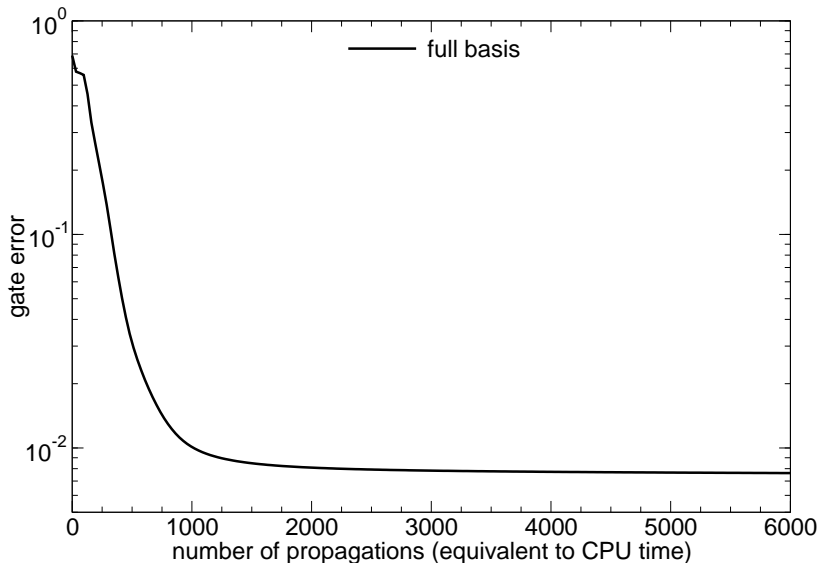
$$\frac{\partial \hat{\rho}}{\partial t} = \frac{i}{\hbar} [\hat{\mathbf{H}}_{\text{eff}}, \hat{\rho}] + \mathcal{L}_D^{(1)}(\hat{\rho}) + \mathcal{L}_D^{(2)}(\hat{\rho})$$

$$\text{with } \mathcal{L}_D^{(q)}(\hat{\rho}) = \gamma_q \sum_{i=1}^{N-1} iD \left[|i-1\rangle\langle i|_q \right] \hat{\rho} + \gamma_{\phi,q} \sum_{i=0}^{N-1} \sqrt{i}D \left[|i\rangle\langle i|_q \right] \hat{\rho}$$

$$\text{with } D \left[\hat{\mathbf{A}} \right] \hat{\rho} = \hat{\mathbf{A}}\hat{\rho}\hat{\mathbf{A}}^\dagger - \frac{1}{2} \left(\hat{\mathbf{A}}^\dagger \hat{\mathbf{A}}\hat{\rho} + \hat{\rho}\hat{\mathbf{A}}^\dagger \hat{\mathbf{A}} \right)$$

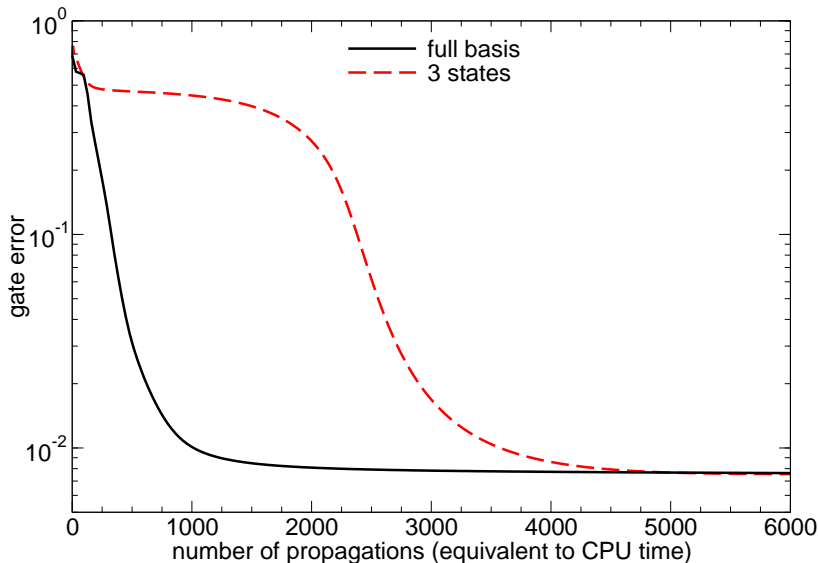
decay time $T_1 = 38.0 \mu\text{s}, 32.0 \mu\text{s}$; dephasing time $T_2^* = 29.5 \mu\text{s}, 16.0 \mu\text{s}$

Optimization of a Transmon Gate – CPU Time



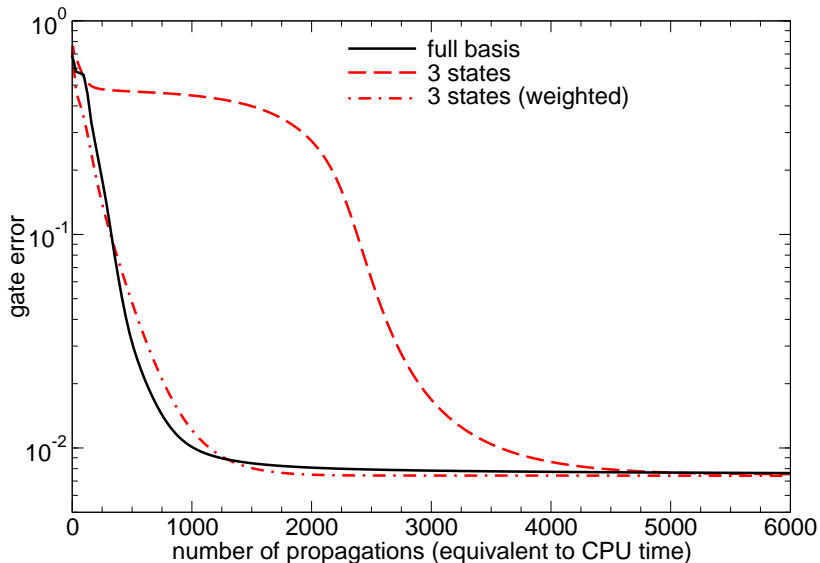
⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP

Optimization of a Transmon Gate – CPU Time



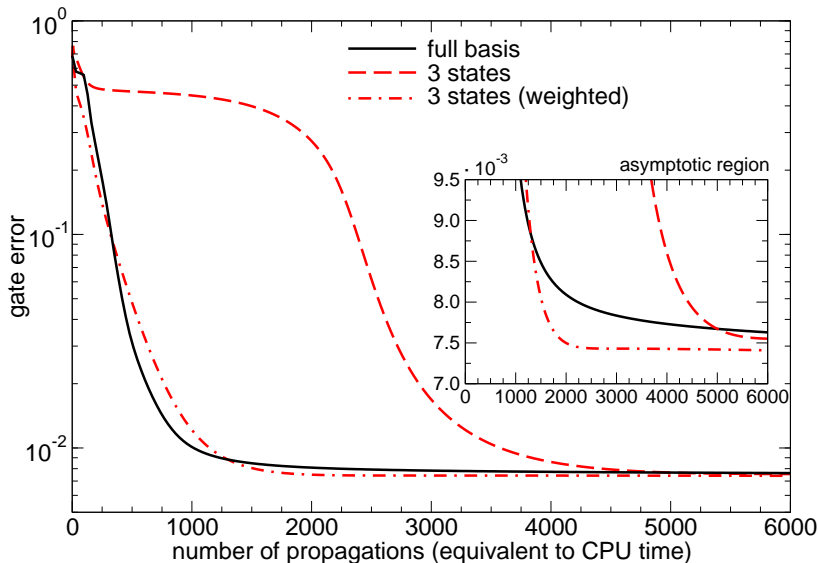
⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP

Optimization of a Transmon Gate – CPU Time



⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP

Optimization of a Transmon Gate – CPU Time



⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP

Using Pure States Only

$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix}, \quad \hat{\rho}_2 = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}, \quad \hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

Using Pure States Only

$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix}, \quad \hat{\rho}_2 = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}, \quad \hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

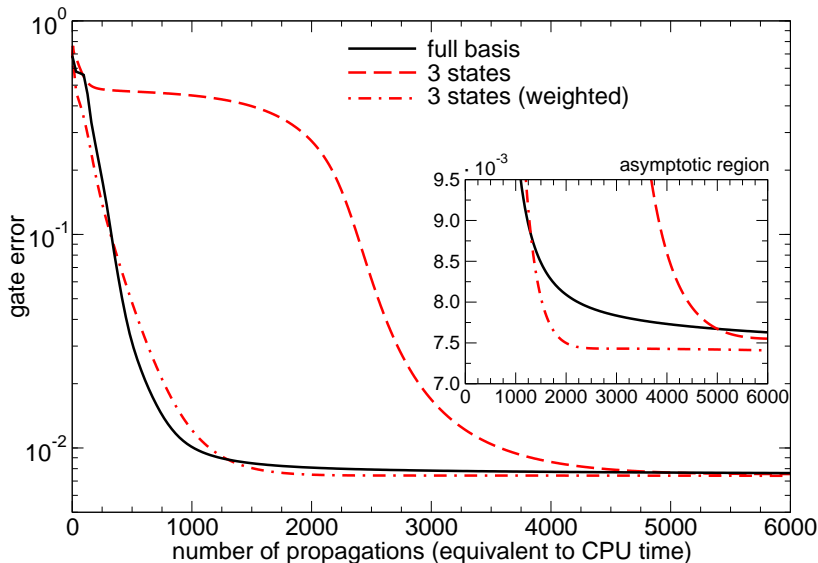
↓

⇓

$$\hat{\rho}_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \quad \hat{\rho}_2 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix},$$

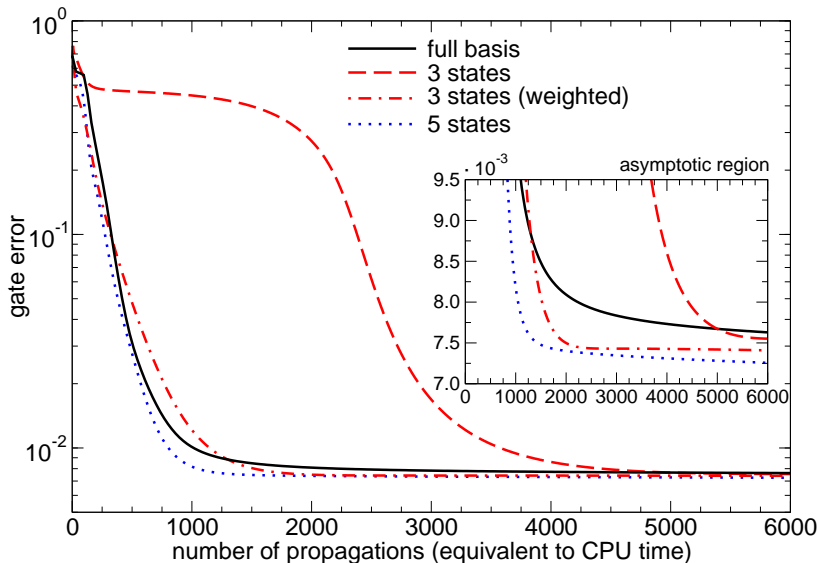
$$\hat{\rho}_3 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \quad \hat{\rho}_4 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad \hat{\rho}_5 = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}$$

Optimization of a Transmon Gate – CPU Time



⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP

Optimization of a Transmon Gate – CPU Time



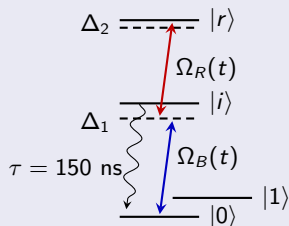
⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP

Example 2

Optimization for Robustness under System Fluctuations for a Rydberg Gate

Two Trapped Neutral Atoms

Single-qubit Hamiltonian

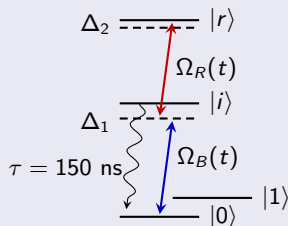


In the RWA:

$$\hat{\mathbf{H}}_{1q} = \begin{pmatrix} 0 & 0 & \frac{1}{2}\Omega_R(t) & 0 \\ 0 & E1 & 0 & 0 \\ \frac{1}{2}\Omega_R(t) & 0 & \Delta_1 & \frac{1}{2}\Omega_B(t) \\ 0 & 0 & \frac{1}{2}\Omega_B(t) & 0 \end{pmatrix}$$

Two Trapped Neutral Atoms

Single-qubit Hamiltonian



In the RWA:

$$\hat{H}_{1q} = \begin{pmatrix} 0 & 0 & \frac{1}{2}\Omega_R(t) & 0 \\ 0 & E1 & 0 & 0 \\ \frac{1}{2}\Omega_R(t) & 0 & \Delta_1 & \frac{1}{2}\Omega_B(t) \\ 0 & 0 & \frac{1}{2}\Omega_B(t) & 0 \end{pmatrix}$$

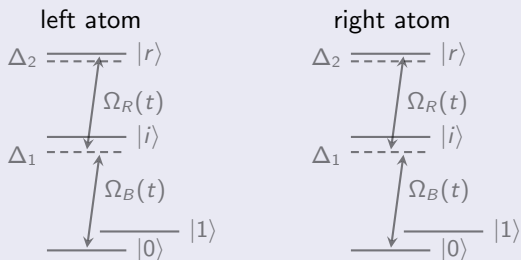
Two-qubit Hamiltonian

$$\hat{H}_{2q} = \hat{H}_{1q} \otimes \mathbb{1} + \mathbb{1} \otimes \hat{H}_{1q} - U |rr\rangle\langle rr|$$

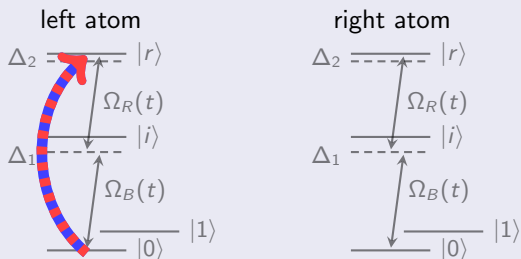
Dipole-dipole interaction when both atoms in Rydberg state.

Only diagonal gates!

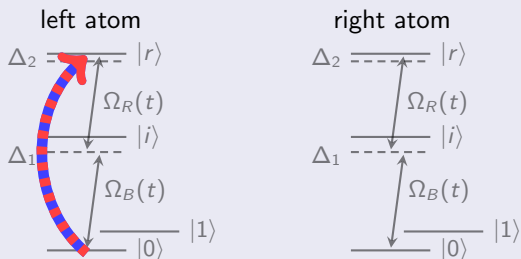
Two trapped Rydberg atoms



Two trapped Rydberg atoms



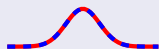
Two trapped Rydberg atoms



Rabi-pulses in three-level systems

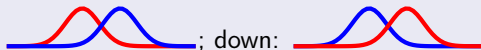
- Option 1:

two-photon pulse, adiabatic elimination of level $|i\rangle$.



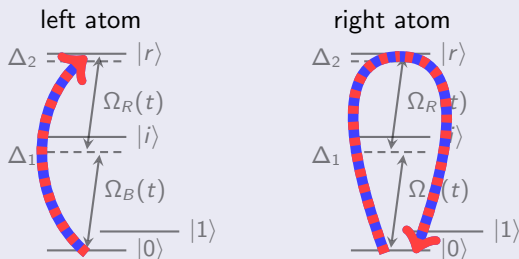
- Option 2:

STIRAP. up:



⇒ combine both options

Two trapped Rydberg atoms



Rabi-pulses in three-level systems

- Option 1:

two-photon pulse, adiabatic elimination of level $|i\rangle$.



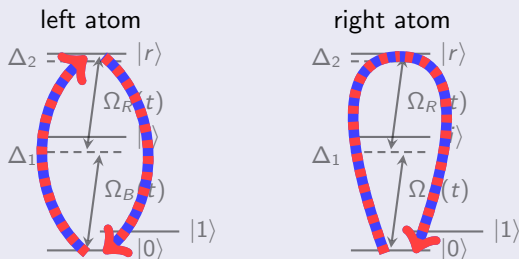
- Option 2:

STIRAP. up:



⇒ combine both options

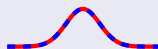
Two trapped Rydberg atoms



Rabi-pulses in three-level systems

- Option 1:

two-photon pulse, adiabatic elimination of level $|i\rangle$.



- Option 2:

STIRAP. up:

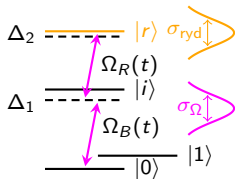


⇒ combine both options

Experimental Fluctuations

Fluctuations:

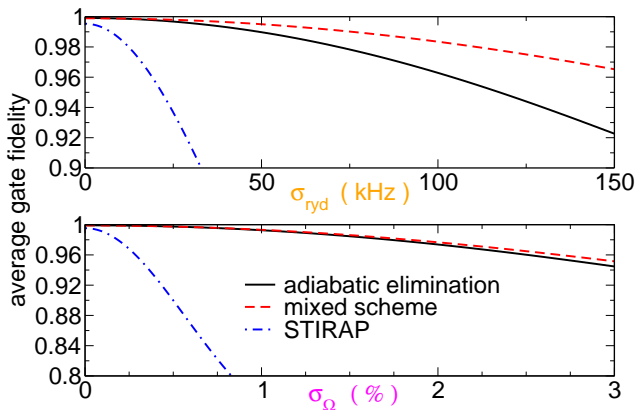
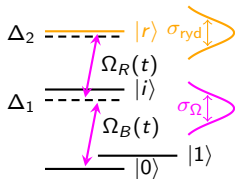
- Rydberg level (width σ_{ryd}) – external fields
- Pulse amplitude (width σ_{Ω})



Experimental Fluctuations

Fluctuations:

- Rydberg level (width σ_{ryd}) – external fields
- Pulse amplitude (width σ_{Ω})



⇒ Goerz, Halperin, Aytac, Koch, Whaley. arXiv:1401.1858.

How can OCT improve the robustness?

Idea: Ensemble Optimization

Generate set of N Hamiltonians, sampling parameter fluctuations.
Optimize for all Hamiltonians *simultaneously*.

Idea: Ensemble Optimization

Generate set of N Hamiltonians, sampling parameter fluctuations.
Optimize for all Hamiltonians *simultaneously*.

Optimization Functional

$$J = 1 - \sum_{n=1}^N \sum_{i=1}^m \frac{w_{i,n}}{\text{tr}[\hat{\rho}_i^2(0)]} \Re \left\{ \text{tr} \left[\hat{\mathbf{O}} \hat{\rho}_i(0) \hat{\mathbf{O}}^\dagger \hat{\rho}_{i,n}(T) \right] \right\}$$

- N : number of realizations $\hat{\mathbf{H}}_n$: $N = 24$
- m : number of states, for each realizations
- $m \times N$ density matrix propagations!

$$\frac{\partial}{\partial t} \hat{\rho}_{i,n}(t) = -i\hbar[\hat{\mathbf{H}}_n(t), \hat{\rho}_{i,n}(t)] + \mathcal{L}_D(\hat{\rho}_{i,n}(t))$$

Idea: Ensemble Optimization

Generate set of N Hamiltonians, sampling parameter fluctuations.
Optimize for all Hamiltonians *simultaneously*.

Optimization Functional

$$J = 1 - \sum_{n=1}^N \sum_{i=1}^m \frac{w_{i,n}}{\text{tr}[\hat{\rho}_i^2(0)]} \Re \left\{ \text{tr} \left[\hat{\mathbf{O}} \hat{\rho}_i(0) \hat{\mathbf{O}}^\dagger \hat{\rho}_{i,n}(T) \right] \right\}$$

- N : number of realizations $\hat{\mathbf{H}}_n$: $N = 24$
- m : number of states, for each realizations ... $m = 3$?
- $m \times N$ density matrix propagations!

$$\frac{\partial}{\partial t} \hat{\rho}_{i,n}(t) = -i\hbar[\hat{\mathbf{H}}_n(t), \hat{\rho}_{i,n}(t)] + \mathcal{L}_D(\hat{\rho}_{i,n}(t))$$

Idea: Ensemble Optimization

Generate set of N Hamiltonians, sampling parameter fluctuations.
Optimize for all Hamiltonians *simultaneously*.

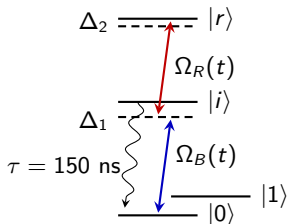
Optimization Functional

$$J = 1 - \sum_{n=1}^N \sum_{i=1}^m \frac{w_{i,n}}{\text{tr}[\hat{\rho}_i^2(0)]} \Re \left\{ \text{tr} \left[\hat{\mathbf{O}} \hat{\rho}_i(0) \hat{\mathbf{O}}^\dagger \hat{\rho}_{i,n}(T) \right] \right\}$$

- N : number of realizations $\hat{\mathbf{H}}_n$: $N = 24$
- m : number of states, for each realizations ... $m = 3$? Not 16!
- $m \times N$ density matrix propagations!

$$\frac{\partial}{\partial t} \hat{\rho}_{i,n}(t) = -i\hbar[\hat{\mathbf{H}}_n(t), \hat{\rho}_{i,n}(t)] + \mathcal{L}_D(\hat{\rho}_{i,n}(t))$$

Diagonal Gates

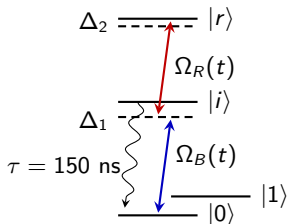


no coupling between $|0\rangle, |1\rangle$

$$\hat{\mathbf{U}} = \text{diag}(e^{i\phi_{00}}, e^{i\phi_{01}}, e^{i\phi_{10}}, e^{i\phi_{11}})$$

only diagonal gates are possible

Diagonal Gates



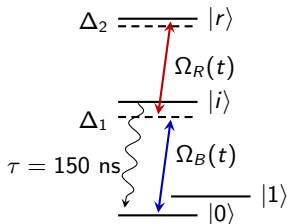
no coupling between $|0\rangle$, $|1\rangle$

$$\hat{\mathbf{U}} = \text{diag}(e^{i\phi_{00}}, e^{i\phi_{01}}, e^{i\phi_{10}}, e^{i\phi_{11}})$$

only diagonal gates are possible

$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix}, \quad \hat{\rho}_2 = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}, \quad \hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

Diagonal Gates



no coupling between $|0\rangle$, $|1\rangle$

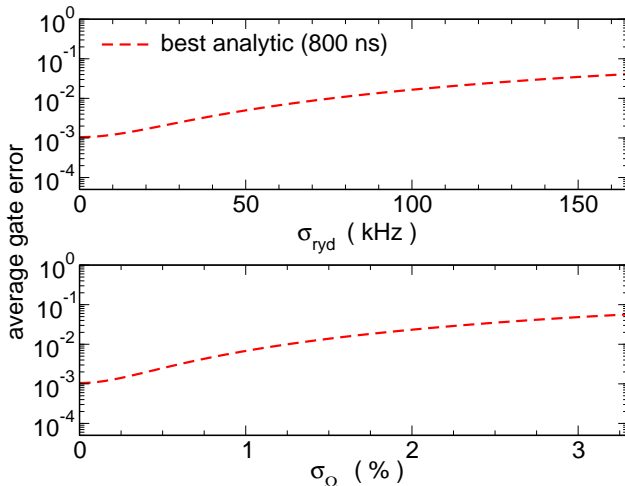
$$\hat{\mathbf{U}} = \text{diag}(e^{i\phi_{00}}, e^{i\phi_{01}}, e^{i\phi_{10}}, e^{i\phi_{11}})$$

only diagonal gates are possible

~~$$\hat{\rho}_1 = \frac{1}{20} \begin{pmatrix} 8 & 0 & 0 & 0 \\ 0 & 6 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix}, \hat{\rho}_2 = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}, \hat{\rho}_3 = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$~~

Robustness of OCT Pulses

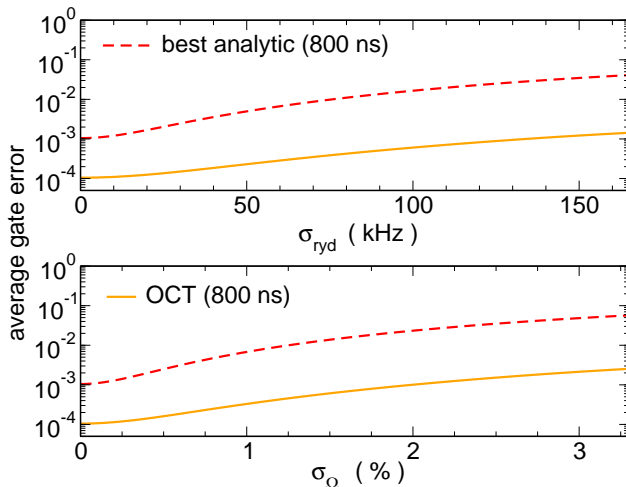
Again: fluctuation of Rydberg level (σ_{Ryd}) and pulse amplitude (σ_{Ω}).



⇒ Goerz, Halperin, Aytac, Koch, Whaley. arXiv:1401.1858.

Robustness of OCT Pulses

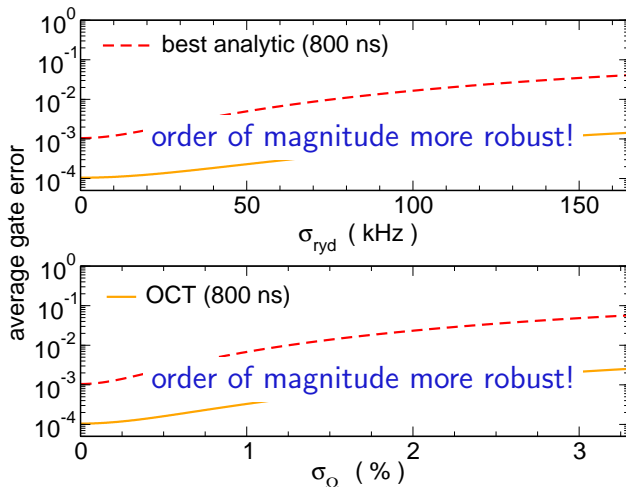
Again: fluctuation of Rydberg level (σ_{Ryd}) and pulse amplitude (σ_{Ω}).



⇒ Goerz, Halperin, Aytac, Koch, Whaley. arXiv:1401.1858.

Robustness of OCT Pulses

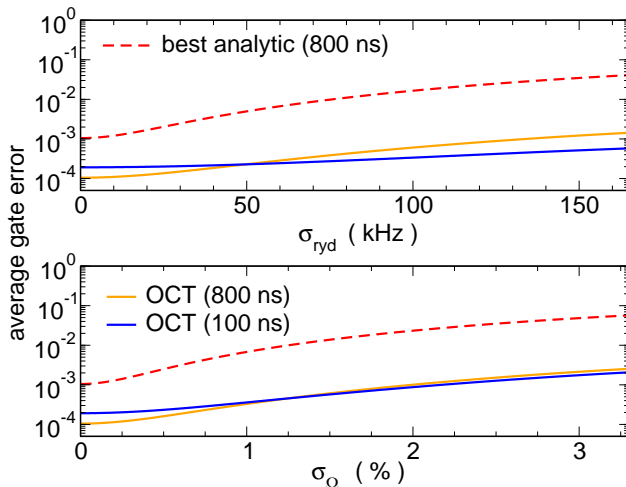
Again: fluctuation of Rydberg level (σ_{Ryd}) and pulse amplitude (σ_{Ω}).



⇒ Goerz, Halperin, Aytac, Koch, Whaley. arXiv:1401.1858.

Robustness of OCT Pulses

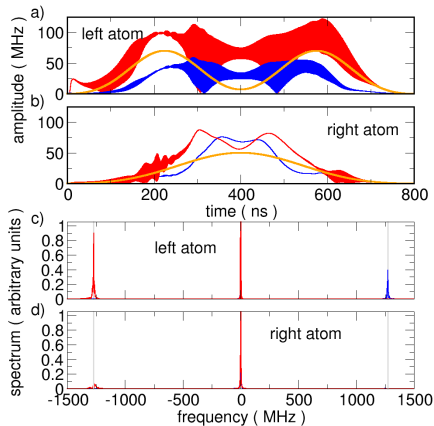
Again: fluctuation of Rydberg level (σ_{Ryd}) and pulse amplitude (σ_{Ω}).



⇒ Goerz, Halperin, Aytac, Koch, Whaley. arXiv:1401.1858.

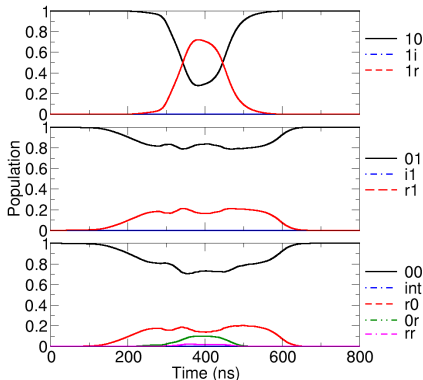
Optimized Pulses & Dynamics

Pulses and Spectra



- STIRAP-like population transfer
- interference between pathways

Population Dynamics



- no population in decaying intermediary states

- Efficiently optimizing for robustness w.r.t. dissipation:
A set of three density matrices is sufficient
(independent of dimension of Hilbert space!)
 - one to check dynamical map on subspace
 - one to check the basis
 - one to check the phases

Further reduction possible for restricted systems!

- Efficiently optimizing for robustness w.r.t. dissipation:
A set of three density matrices is sufficient
(independent of dimension of Hilbert space!)
 - one to check dynamical map on subspace
 - one to check the basis
 - one to check the phases

Further reduction possible for restricted systems!

- Optimize for Robustness w.r.t. fluctuations in experimental parameters: ensemble optimization
⇒ Example: highly robust Rydberg gates

- Efficiently optimizing for robustness w.r.t. dissipation:
A set of three density matrices is sufficient
(independent of dimension of Hilbert space!)
 - one to check dynamical map on subspace
 - one to check the basis
 - one to check the phases

Further reduction possible for restricted systems!

- Optimize for Robustness w.r.t. fluctuations in experimental parameters: ensemble optimization
⇒ Example: highly robust Rydberg gates

Reference:

- ⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP
- ⇒ Goerz, Halperin, Aytac, Koch, Whaley. arXiv:1401.1858.

Acknowledgments

Kassel:



- Christiane Koch
- Daniel Reich

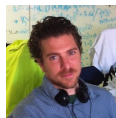
Berkeley:



■ Birgitta Whaley



■ Eli Halperin



■ Jon Aytac

Acknowledgments

Kassel:



- Christiane Koch
- Daniel Reich

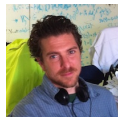
Berkeley:



Birgitta Whaley



Eli Halperin



Jon Aytac

Thank you

- Efficiently optimizing for robustness w.r.t. dissipation:
A set of three density matrices is sufficient
(independent of dimension of Hilbert space!)
 - one to check dynamical map on subspace
 - one to check the basis
 - one to check the phases

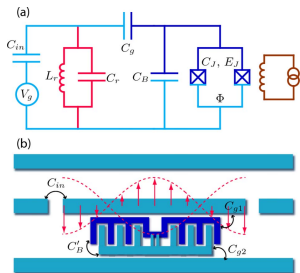
Further reduction possible for restricted systems!

- Optimize for Robustness w.r.t. fluctuations in experimental parameters: ensemble optimization
⇒ Example: highly robust Rydberg gates

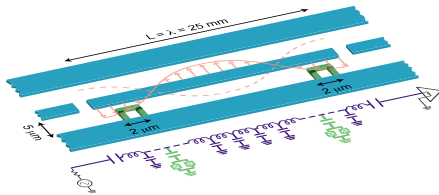
Reference:

- ⇒ Goerz, Reich, Koch. arXiv:1312.0111. In press: NJP
- ⇒ Goerz, Halperin, Aytac, Koch, Whaley. arXiv:1401.1858.

Two Coupled Transmon Qubits

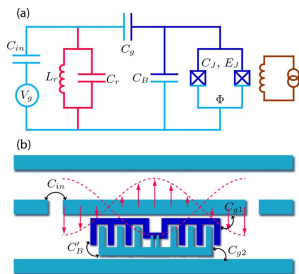


J. Koch et al. PRA 76, 042319 (2007)

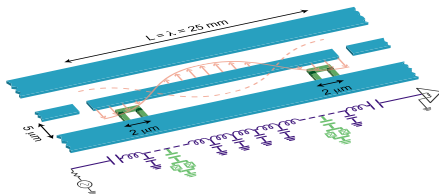


A. Blais et al. PRA 75, 032329 (2007)

Two Coupled Transmon Qubits



J. Koch et al. PRA 76, 042319 (2007)



A. Blais et al. PRA 75, 032329 (2007)

Full Hamiltonian

$$\hat{H} = \underbrace{\omega_c \hat{a}^\dagger \hat{a}}_{(1)} + \underbrace{\omega_1 \hat{b}_1^\dagger \hat{b}_1 + \omega_2 \hat{b}_2^\dagger \hat{b}_2}_{(2)} - \underbrace{\frac{1}{2} (\alpha_1 \hat{b}_1^\dagger \hat{b}_1^\dagger \hat{b}_1 \hat{b}_1 + \alpha_2 \hat{b}_2^\dagger \hat{b}_2^\dagger \hat{b}_2 \hat{b}_2)}_{(3)} + \underbrace{g_1 (\hat{b}_1^\dagger \hat{a} + \hat{b}_1 \hat{a}^\dagger) + g_2 (\hat{b}_2^\dagger \hat{a} + \hat{b}_2 \hat{a}^\dagger)}_{(4)} + \underbrace{\epsilon^*(t) \hat{a} + \epsilon(t) \hat{a}^\dagger}_{(5)}$$

$$\begin{aligned}\hat{\mathbf{H}}_{\text{eff}} = & \sum_{q=1,2} \sum_{i=0}^{N_q-1} (\omega_i^{(q)} + \chi_i^{(q)}) \hat{\mathbf{n}}_i^{(q)} + \sum_{q=1,2} \sum_{i=0}^{N_q-1} g_i^{\text{eff}(q)} \epsilon(t) (\hat{\mathbf{c}}_i^{+(q)} + \hat{\mathbf{c}}_i^{- (q)}) \\ & + \sum_{ij} J_{ij}^{\text{eff}} (\hat{\mathbf{c}}_i^{- (1)} \hat{\mathbf{c}}_j^{+(2)} + \hat{\mathbf{c}}_i^{+(1)} \hat{\mathbf{c}}_j^{- (2)}).\end{aligned}$$

$$\hat{H}_{\text{eff}} = \sum_{q=1,2} \sum_{i=0}^{N_q-1} (\omega_i^{(q)} + \chi_i^{(q)}) \hat{\Pi}_i^{(q)} + \sum_{q=1,2} \sum_{i=0}^{N_q-1} g_i^{\text{eff}(q)} \epsilon(t) (\hat{\mathbf{C}}_i^{+(q)} + \hat{\mathbf{C}}_i^{- (q)}) \\ + \sum_{ij} J_{ij}^{\text{eff}} (\hat{\mathbf{C}}_i^{- (1)} \hat{\mathbf{C}}_j^{+(2)} + \hat{\mathbf{C}}_i^{+(1)} \hat{\mathbf{C}}_j^{- (2)}).$$

with

- $\omega_i^{(q)} = i\omega_q - \frac{1}{2}(i^2 - i)\alpha_q, \quad g_i^{(q)} = \sqrt{i}g_q$
- $\hat{\Pi}_i^{(q)} = |i\rangle\langle i|_q, \quad \hat{\mathbf{C}}_i^{+(q)} = |i\rangle\langle i-1|_q$
- $\chi_i^{(q)} = \frac{(g_i^{(q)})^2}{(\omega_i^{(q)} - \omega_{i-1}^{(q)} - \omega_c)}$
- $g_i^{\text{eff}(q)} = \frac{g_i^{(q)}}{(\omega_i^{(q)} - \omega_{i-1}^{(q)} - \omega_c)}$
- $J_{ij}^{\text{eff}} = \frac{1}{2}g_i^{\text{eff}(1)}g_j^{(2)} + \frac{1}{2}g_j^{\text{eff}(2)}g_i^{(1)}$

qubit frequency ω_1	4.3796 GHz
qubit frequency ω_2	4.6137 GHz
drive frequency ω_d	4.4985 GHz
anharmonicity α_1	-239.3 MHz
anharmonicity α_2	-242.8 MHz
effective qubit-qubit coupling J	-2.3 MHz
qubit 1,2 decay time T_1	38.0 μ s, 32.0 μ s
qubit 1,2 dephasing time T_2^*	29.5 μ s, 16.0 μ s

Effective Hamiltonian

$$\hat{H}_{\text{eff}} = \sum_{ijq} \left((\omega_i^{(q)} + \chi_i^{(q)}) \hat{n}_i^{(q)} + g_i^{\text{eff}(q)} \epsilon(t) (\hat{C}_i^{+(q)} + \hat{C}_i^{-(q)}) + J_{ij}^{\text{eff}} (\hat{C}_i^{-(1)} \hat{C}_j^{+(2)} + \text{c.c.}) \right)$$

Master Equation

$$\mathcal{L}_D(\hat{\rho}) = \sum_{q=1,2} \left(\gamma_q \sum_{i=1}^{N-1} iD \left[|i-1\rangle\langle i|_q \right] \hat{\rho} + \gamma_{\phi,q} \sum_{i=0}^{N-1} \sqrt{i} iD \left[|i\rangle\langle i|_q \right] \hat{\rho} \right),$$

$$\text{with } D[\hat{A}] \hat{\rho} = \hat{A} \hat{\rho} \hat{A}^\dagger - \frac{1}{2} (\hat{A}^\dagger \hat{A} \hat{\rho} + \hat{\rho} \hat{A}^\dagger \hat{A})$$

qubit frequency ω_1	4.3796 GHz
qubit frequency ω_2	4.6137 GHz
drive frequency ω_d	4.4985 GHz
anharmonicity α_1	-239.3 MHz
anharmonicity α_2	-242.8 MHz
effective qubit-qubit coupling J	-2.3 MHz
qubit 1,2 decay time T_1	38.0 μ s, 32.0 μ s
qubit 1,2 dephasing time T_2^*	29.5 μ s, 16.0 μ s

- Near resonance of α_1 with $\omega_1 - \omega_2$

Effective Hamiltonian

$$\hat{H}_{\text{eff}} = \sum_{ijq} \left((\omega_i^{(q)} + \chi_i^{(q)}) \hat{n}_i^{(q)} + g_i^{\text{eff}(q)} \epsilon(t) (\hat{C}_i^{+(q)} + \hat{C}_i^{-(q)}) + J_{ij}^{\text{eff}} (\hat{C}_i^{-(1)} \hat{C}_j^{+(2)} + \text{c.c.}) \right)$$

Master Equation

$$\mathcal{L}_D(\hat{\rho}) = \sum_{q=1,2} \left(\gamma_q \sum_{i=1}^{N-1} iD \left[|i-1\rangle\langle i|_q \right] \hat{\rho} + \gamma_{\phi,q} \sum_{i=0}^{N-1} \sqrt{i} iD \left[|i\rangle\langle i|_q \right] \hat{\rho} \right),$$

$$\text{with } D[\hat{A}] \hat{\rho} = \hat{A} \hat{\rho} \hat{A}^\dagger - \frac{1}{2} (\hat{A}^\dagger \hat{A} \hat{\rho} + \hat{\rho} \hat{A}^\dagger \hat{A})$$

qubit frequency ω_1	4.3796 GHz
qubit frequency ω_2	4.6137 GHz
drive frequency ω_d	4.4985 GHz
anharmonicity α_1	-239.3 MHz
anharmonicity α_2	-242.8 MHz
effective qubit-qubit coupling J	-2.3 MHz
qubit 1,2 decay time T_1	38.0 μ s, 32.0 μ s
qubit 1,2 dephasing time T_2^*	29.5 μ s, 16.0 μ s

- Near resonance of α_1 with $\omega_1 - \omega_2$
- single frequency drive centered between two qubits

Effective Hamiltonian

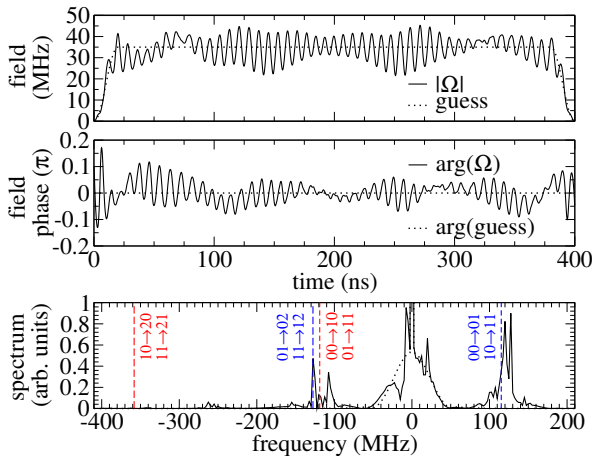
$$\hat{H}_{\text{eff}} = \sum_{ijq} \left((\omega_i^{(q)} + \chi_i^{(q)}) \hat{n}_i^{(q)} + g_i^{\text{eff}(q)} \epsilon(t) (\hat{C}_i^{+(q)} + \hat{C}_i^{-(q)}) + J_{ij}^{\text{eff}} (\hat{C}_i^{-(1)} \hat{C}_j^{+(2)} + \text{c.c.}) \right)$$

Master Equation

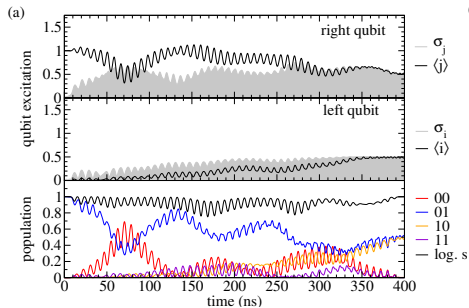
$$\mathcal{L}_D(\hat{\rho}) = \sum_{q=1,2} \left(\gamma_q \sum_{i=1}^{N-1} iD \left[|i-1\rangle\langle i|_q \right] \hat{\rho} + \gamma_{\phi,q} \sum_{i=0}^{N-1} \sqrt{i} iD \left[|i\rangle\langle i|_q \right] \hat{\rho} \right),$$

$$\text{with } D[\hat{A}] \hat{\rho} = \hat{A} \hat{\rho} \hat{A}^\dagger - \frac{1}{2} (\hat{A}^\dagger \hat{A} \hat{\rho} + \hat{\rho} \hat{A}^\dagger \hat{A})$$

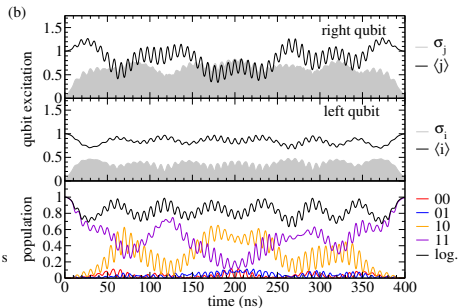
Transmon Optimized Pulse



Transmon Population Dynamics



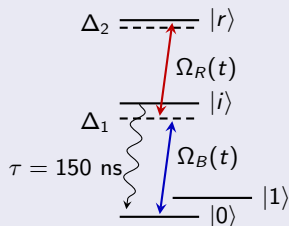
$$\Psi(t = 0) = |01\rangle$$



$$\Psi(t = 0) = |11\rangle$$

Analytical Gate Scheme

Single-qubit Hamiltonian



In the RWA:

$$\hat{H}_{1q} = \begin{pmatrix} 0 & 0 & \frac{1}{2}\Omega_R(t) & 0 \\ 0 & E1 & 0 & 0 \\ \frac{1}{2}\Omega_R(t) & 0 & \Delta_1 & \frac{1}{2}\Omega_B(t) \\ 0 & 0 & \frac{1}{2}\Omega_B(t) & 0 \end{pmatrix}$$

Gate scheme (in blockade regime)

	π (left)		2π (right)		π (left)	
$ 00\rangle$	\rightarrow	$i r0\rangle$	\equiv	$i r0\rangle$	\rightarrow	$- 00\rangle$
$ 10\rangle$	\rightarrow	$ 10\rangle$	\rightarrow	$- 10\rangle$	\rightarrow	$- 10\rangle$
$ 01\rangle$	\rightarrow	$i r1\rangle$	\rightarrow	$i r1\rangle$	\rightarrow	$- 01\rangle$
$ 11\rangle$	\rightarrow	$ 11\rangle$	\rightarrow	$ 11\rangle$	\rightarrow	$ 11\rangle$